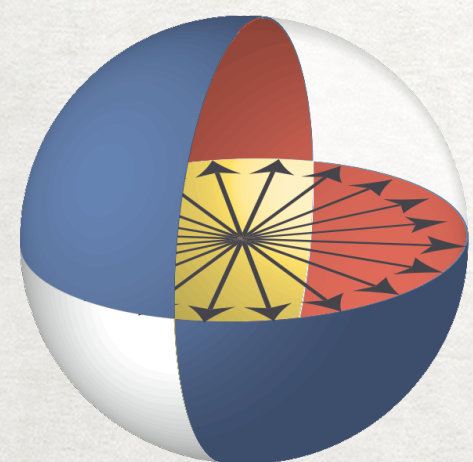


Local Hamiltonians with No Low-Energy Stabilizer States

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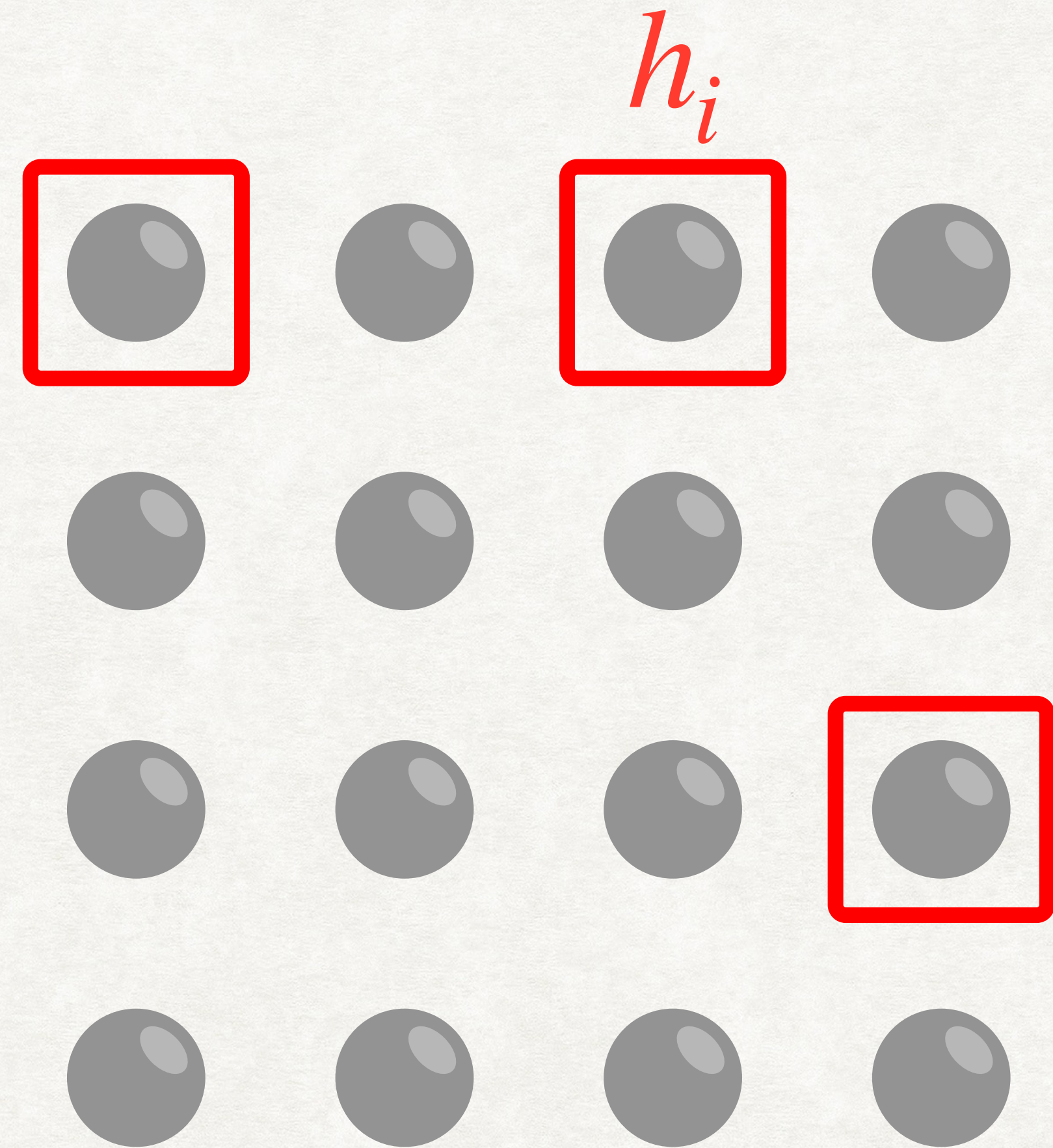
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Local Hamiltonians

k -local interaction term: h_i PSD with $\|h_i\| \leq 1$



Local Hamiltonians

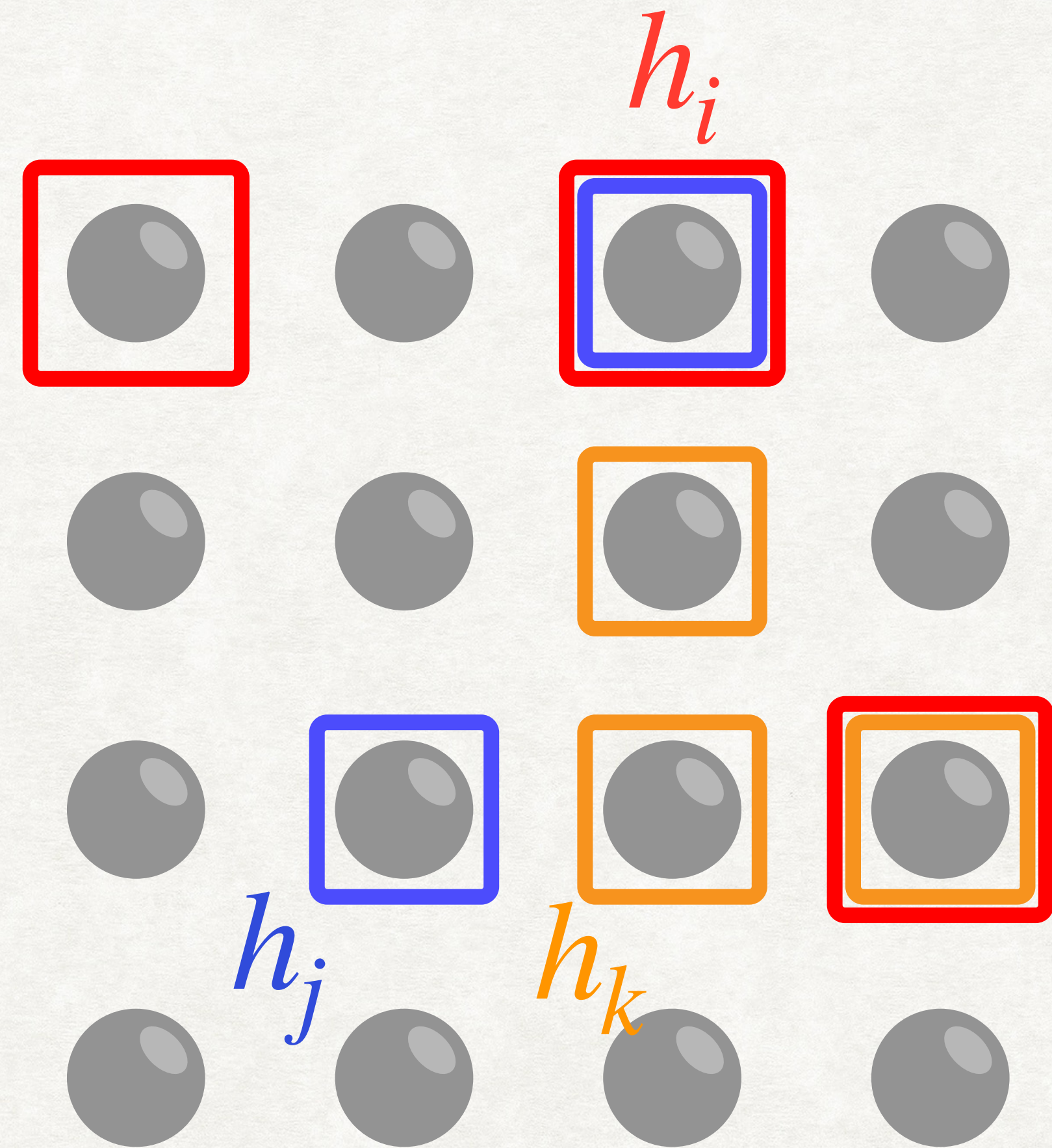
k -local interaction term: h_i PSD with $\|h_i\| \leq 1$

k -local Hamiltonian: $m = \text{poly}(n)$ k -local terms

$$H = \frac{1}{m} \sum_{i=1}^m h_i \otimes I_{2^{n-k}}$$

Ground-state energy: $E_{gs} = \min_{|\psi\rangle} \langle \psi | H | \psi \rangle$

Can we approximate E_{gs} to within some error $\delta(n)$ (in BQP)?



Local Hamiltonian Problem

Local Hamiltonian problem (LH): given H , $\epsilon > 0$, $\delta(n)$, decide between

$$(1) E_{gs} \leq \epsilon \quad \text{or} \quad (2) E_{gs} > \epsilon + \delta(n).$$

For $\delta(n) = \frac{1}{\text{poly}(n)}$, LH is QMA-complete [KSV02].

\Rightarrow Can't approximate E_{gs} with error $\delta(n) = \frac{1}{\text{poly}(n)}$ unless BQP=QMA.

qPCP Conjecture

Can we approximate E_{gs} with constant precision?

Quantum PCP Conjecture (qPCP): LH is QMA-hard with $\delta(n) = \Omega(1)$.

By classical PCP Theorem LH- $\Omega(1)$ is at least NP-hard.

What does a "qPCP Hamiltonian" look like?

If $C \subsetneq \text{QMA} = \text{QPCP}$ is conjectured \Rightarrow cannot estimate $E_{gs} \pm \epsilon$ in C .

\Rightarrow No ground state should have an energy estimation algorithm in C .

What do *low-energy states* look like?

If $C \subsetneq \text{QMA} = \text{QPCP}$ is conjectured \Rightarrow cannot estimate $E_{gs} \pm \epsilon$ in C .

\Rightarrow No *low-energy state* should have an energy estimation algorithm in C .

[Low-energy = small constant energy above E_{gs}]

If qPCP Conjecture is true, $\exists H$ whose low-energy states cannot admit energy approximation algorithms.

Some ways to estimate energy

1. Trivial (i.e. Low-depth circuit) states — NP via light cone argument
2. "Sampleable states" — MA via dequantizing QSVT [GL22]
3. Stabilizer states — NP via stabilizer generators
4. Matrix product states — NP via tensor contraction
5. ...

Some ways to estimate energy

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A "qPCP" Hamiltonian can't have any of these in its low-energy space.

Low-energy space implications

Hamiltonians that should exist if qPCP is true...

1. No low-energy trivial states — **NLTS Theorem [ABN22]**
2. No low-energy “sampleable states” — **NLSS Conjecture [GL22]**
3. No low-energy stabilizer states — **NLCS Theorem (our work)**

Note: Stabilizer states can be efficiently sampleable, so NLSS \Rightarrow NLCS.

No Low-Energy Stabilizer States (NLCS)

Stabilizer state \Leftrightarrow prepared by a Clifford circuit (H, CNOT, S)

H satisfies the ϵ -NLCS property if every stabilizer state has energy $\langle \psi | H | \psi \rangle \geq \epsilon$.*

Intuition: Ground-state energy can't be approximated using stabilizer states.

Theorem. There exists an explicit $\sin^2(\pi/8)$ -NLCS Hamiltonian.

* if $E_{gs} = 0$

Simple NLCS Hamiltonian

Starting point: $H_0 \equiv \frac{1}{n} \sum |1\rangle\langle 1|_i$

Ground state: $|0^n\rangle$ with 0 energy

Main idea: rotate the ground-space into a basis which is highly *non-stabilizer*.

We consider the Y version of the T gate: $D \equiv e^{i\frac{\pi}{8}Y}$

Simple NLCS Hamiltonian

Rotated: $\tilde{H}_0 \equiv D^{\otimes n} H_0 D^{\dagger \otimes n} = \frac{1}{n} \sum D |1\rangle \langle 1|_i D^{\dagger}$

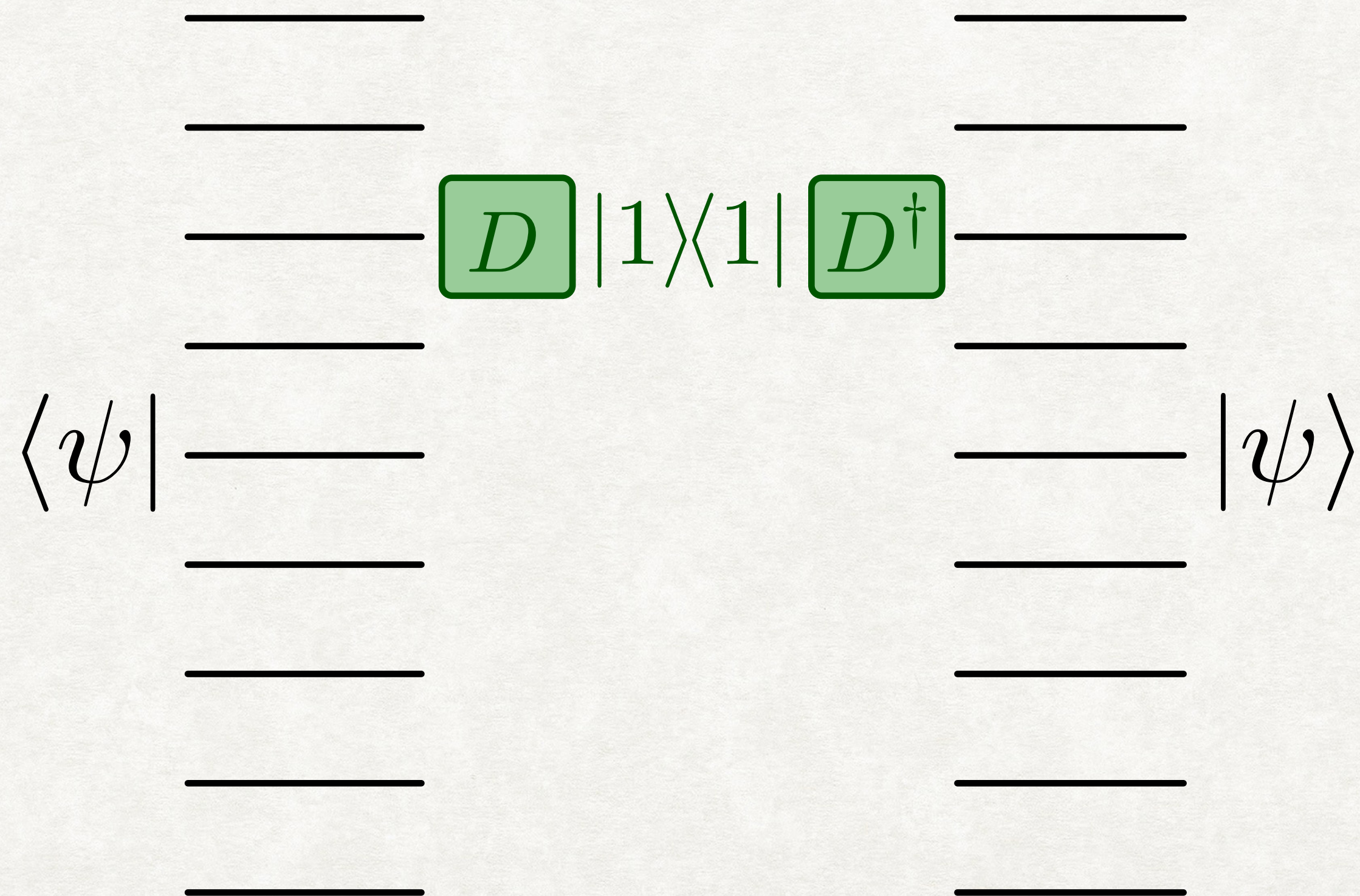
New ground state: $D^{\otimes n} |0^n\rangle$ with 0 energy

Theorem. \tilde{H}_0 is $\sin^2(\pi/8)$ -NLCS.

Simple NLCS Hamiltonian

Theorem. \tilde{H}_0 is $\sin^2(\pi/8)$ -NLCS.

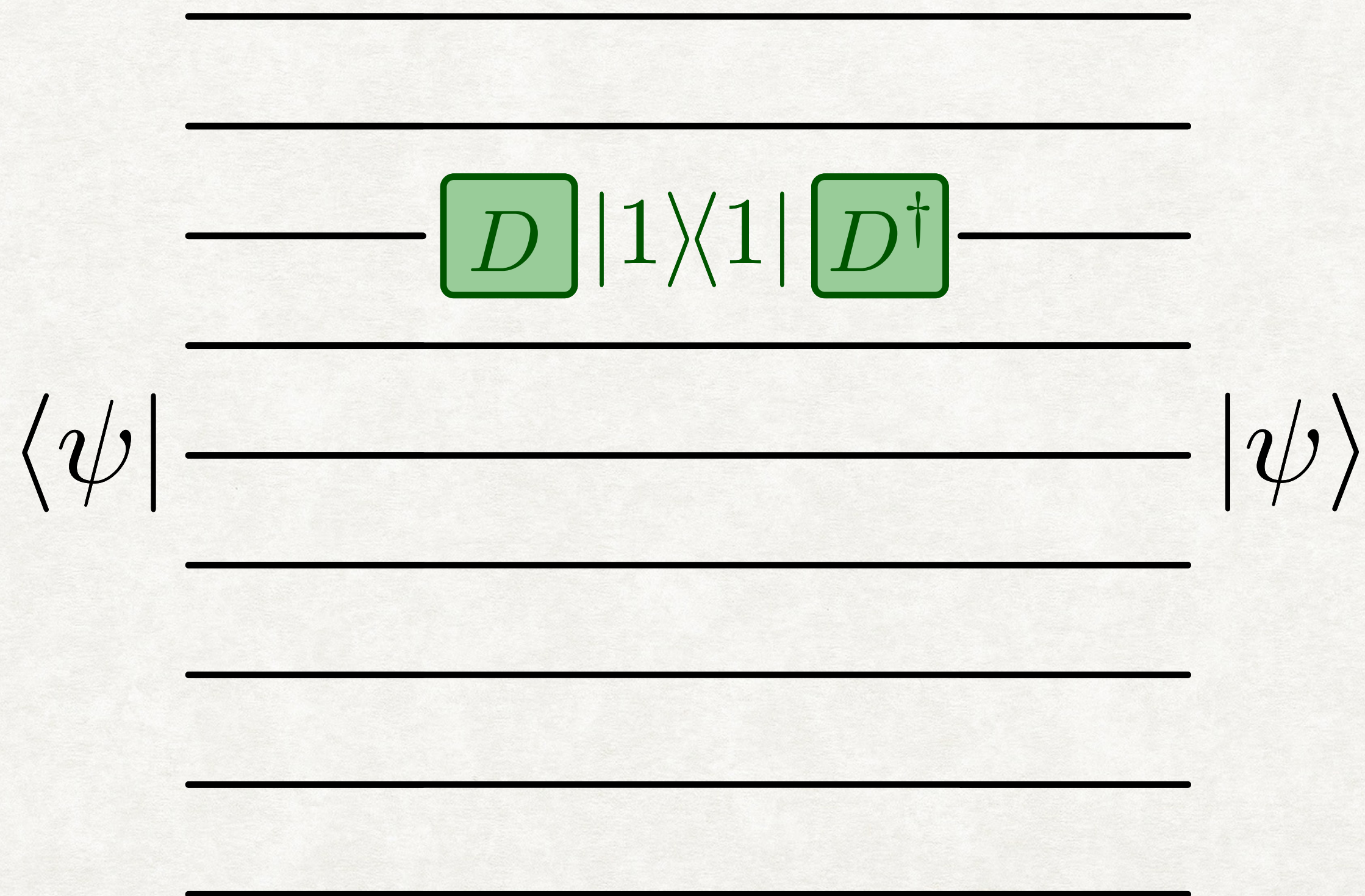
Proof. Consider a single term:



Simple NLCS Hamiltonian

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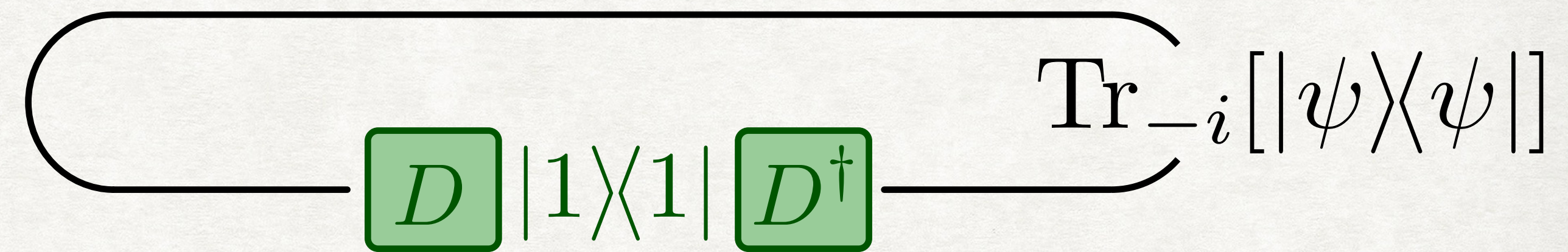
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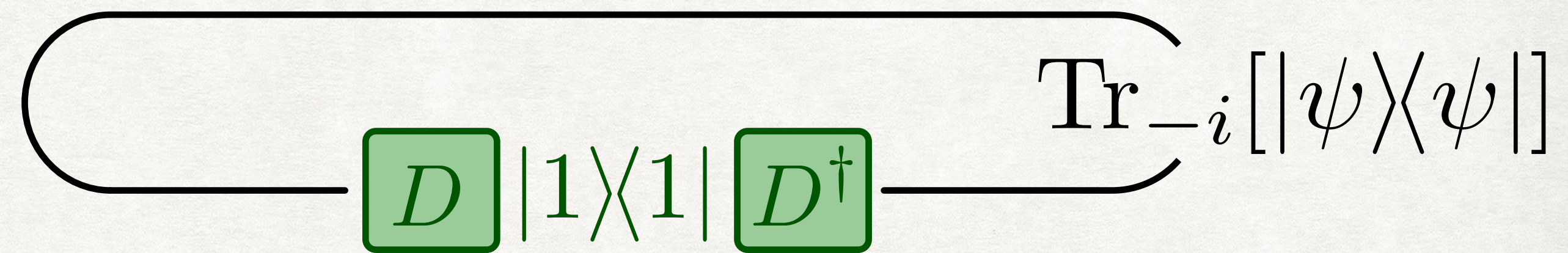


Fact: If $|\psi\rangle$ is an n -qubit stabilizer state then $\psi_i \equiv \text{Tr}_{-i}[|\psi\rangle\langle\psi|]$ is either $\frac{1}{2}I$ or a stabilizer state $|\eta\rangle\langle\eta|$.

Simple NLCS Hamiltonian

Theorem. \tilde{H}_0 is $\sin^2(\pi/8)$ -NLCS.

Proof. Consider a single term:



So:

$$\langle\psi|D|1\rangle\langle 1|_i D^\dagger|\psi\rangle = \frac{1}{2}$$

$$\text{or} = |\langle\eta|D|1\rangle|^2$$

Simple NLCS Hamiltonian

Theorem. \tilde{H}_0 is $\sin^2(\pi/8)$ -NLCS.

Proof.

Direct computation: $|\langle \eta | D | 1 \rangle|^2 \geq \sin^2(\pi/8)$ for all single-qubit stabilizer states.

Since all local terms are lower-bounded by $\sin^2(\pi/8)$ we're done.

Combinations of types

Can we get a *simultaneous* NLTS and NLCS Hamiltonian?

Well... $H_{NLTS} \otimes I + I \otimes H_{NLCS}$, but this isn't a very interesting system.

H_{NLTS} is a CSS Hamiltonian, i.e. it's ground-space corresponds to a stabilizer code.

Can we get NLTS+NLCS directly from rotating by $D = e^{i\frac{\pi}{8}Y}$?

CSS Hamiltonians

Theorem. For CSS Hamiltonians, $\tilde{H} \equiv D^{\otimes n} H D^{\dagger \otimes n}$ is $\sin^2(\pi/8)$ -NLCS.

Corollary. \tilde{H}_{NLTS} satisfies both NLTS and NLCS.

Why? H_{NLTS} is a CSS Hamiltonian, and rotating by a constant-depth circuit preserves NLTS.

CSS Hamiltonians

Theorem. For CSS Hamiltonians, $\tilde{H} \equiv D^{\otimes n} H D^{\dagger \otimes n}$ is $\sin^2(\pi/8)$ -NLCS.

Proof components: Let $|\psi\rangle$ be an n -qubit stabilizer state.

(1) k -local states of $|\psi\rangle$ are mixtures of k -qubit stabilizer states \Rightarrow sufficient to bound local energy terms for k -qubit stabilizer states.

(2) The local terms of \tilde{H} look like $\frac{I - H^{\otimes k}}{2}$ or $\frac{I - (-XHX)^{\otimes k}}{2}$ ($H = \text{Hadamard}$) \Rightarrow local lower bound follows from an upper bound on $|\langle \eta | H^{\otimes k} | \eta \rangle|$.

(3) Main technical lemma: For all k -qubit stabilizer states $\Rightarrow |\langle \eta | H^{\otimes k} | \eta \rangle| \leq \frac{1}{\sqrt{2}}$.
* if k is odd

Future Work

- We showed NLCS+NLTS Hamiltonians exist. What if we relax the stabilizer requirement to Clifford + a single T gate? $\log(n)$ T gates? (No Low-Energy "Almost Clifford" States)

"Almost Clifford" states

Take \tilde{H}_0 to be the rotated zero Hamiltonian.

Conjecture. Suppose $|\psi\rangle$ can be prepared with $\leq \alpha T$ gates. Then

$$\langle \psi | \tilde{H}_0 | \psi \rangle \geq \left(1 - \frac{\alpha}{n} \right) \sin^2(\pi/8).$$

"Almost Clifford" states

Take \tilde{H}_0 to be the rotated zero Hamiltonian.

NLACS Theorem (unpublished). Suppose $|\psi\rangle$ can be prepared with $\leq \alpha T$ gates. Then

$$\langle \psi | \tilde{H}_0 | \psi \rangle \geq \left(1 - \frac{\alpha}{n} \right) \sin^2(\pi/8).$$

Corollary. Low-energy states of \tilde{H}_0 require $n - o(n)$ T gates.

$\Rightarrow E_{gs}$ energy can't be approximated using stabilizer + $\log(n)$ T gate states

Note: \tilde{H}_0 still has an NP witness

Future Work

- ~~We showed NLCS Hamiltonians exist. What if we relax the stabilizer requirement to Clifford + a single T gate? $\log(n)$ T gates?~~
- NLSS Conjecture? Suggestion: CH_0C^\dagger for a low-depth circuit C .
- Witness state “lower-bounds”, i.e. Hamiltonians with no low-energy:
 1. “stabilizer then low-depth” states (rotated stabilizer Hamiltonians have NP witnesses)
 2. states with classical descriptions that can be used to compute k -local reduced density matrices (NLLS Conjecture)
- Complexity lower-bounds for constant-gap LH problem (BQP-hardness, MA-hardness, etc.)

Thanks!

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Simple NLACS Hamiltonian

Rotated zero Hamiltonian: $\tilde{H}_0 \equiv \frac{1}{n} \sum D |1\rangle\langle 1|_i D^\dagger \sim \frac{1}{n} \sum \frac{I - H_i}{2}$

Energy of a single term: $\frac{1}{2} \left(1 - \langle \psi | H_i | \psi \rangle \right)$

If $S \in \text{Stab}(|\psi\rangle)$ acts non-trivially on i , then $\langle \psi | H_i | \psi \rangle \leq \frac{1}{\sqrt{2}}$

If $|\psi\rangle$ can be prepared with α T gates, then $n - \alpha$ qubits are acted on non-trivially by a stabilizer.

